

Quantum gates

How to implement a XOR gate?

a \ b	0	1
0	0	1
1	1	0

but $(a, b) \mapsto a \text{ XOR } b$ not unitary

hence $(a, b) \mapsto (a, a \text{ XOR } b)$

(a, b)	$ 00\rangle$	$ 01\rangle$	$ 10\rangle$	$ 11\rangle$
$(a, a \text{ XOR } b)$	$ 00\rangle$	$ 01\rangle$	$ 11\rangle$	$ 10\rangle$

is permutation \Rightarrow unitary

This operation is called "Controlled Not", "CNOT"

In circuit diagrams, this is denoted by



if we identify $|00\rangle, \dots, |11\rangle$ with e_1, \dots, e_4

we may represent CNOT by the unitary matrix

$$\begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{pmatrix}$$

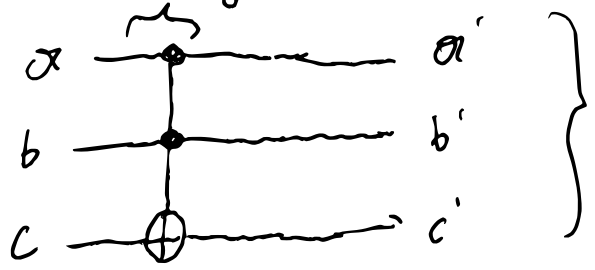
What about "AND"?

$a \text{ AND } b \Leftrightarrow$

$a \backslash b$	0	1
0	0	0
1	0	1

- $(a, b) \mapsto (a, a \text{ AND } b)$ not bijective
- requires three qbits $|a, b, c\rangle \mapsto |a, b, c \oplus (a \text{ AND } b)\rangle$

Toffoli gate



$ abc\rangle$	000	001	010	011	100	101	110	111
	000	001	010	011	100	101	111	110

If input $c = |0\rangle \Rightarrow$ output $c' = |a, b, a \text{ AND } b\rangle$

- extra qbit $|c\rangle$ is called "ancilla" qbit
ancilla = Diener

Deutsch-Jozsa Algorithm

"Problem": Given $f(x_1, \dots, x_n) \in \{0, 1\}$, $x_i \in \{0, 1\}$
determine if

- $f = 0$

- $f = 1$

- f is *balanced*; i.e., exactly half of (x_1, \dots, x_n) lead to $f(x_1, \dots, x_n) = 1$ the other half $\Rightarrow f(x_1, \dots, x_n) = 0$

- We assume the f satisfies one of the three options

- Not useful in practice, but demonstrates that hard problems can be easier on quantum computers

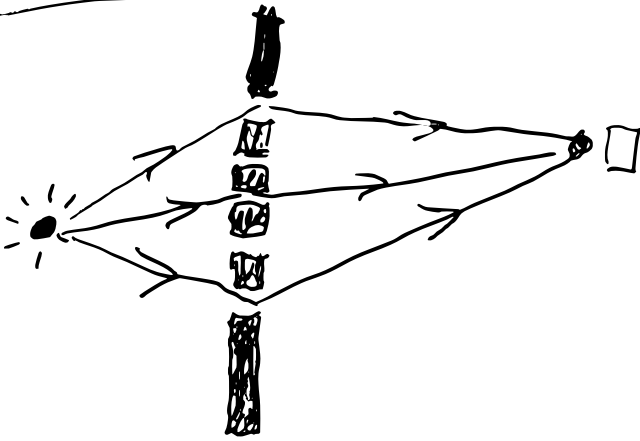
Classical algorithm:

To exclude the *balanced* case, we need to check at least $2^{n-1} + 1$ inputs

$$(x_1, \dots, x_n) \in \{0, 1\}^n$$

\Rightarrow Problem size $N = 2^n$, Runtime $N/2 + 1$

Physical inspiration



Intensity at
detection is determined
by "phase" difference
of light-paths
→ Interference

Number the holes with
 $x \in \{0, 1\}^n$ } 2^n holes

Assume the holes are located such
that the light amplitude is $\approx (-1)^{F(y)}$ at
detector if light travels through hole y .

Light intensity at detector is given
by

$$\approx \sum_{y \in \{0, 1\}^n} (-1)^{F(y)} = 0 \text{ if } F \text{ is balanced}$$
$$= \pm 1 \text{ if } F \text{ is const}$$

Quantum Algorithm

We may identify $(x_1, \dots, x_n) \in \{0, 1\}^n$ with the basis element $|i\rangle$, where $i = x_1 + x_2 \cdot 2 + x_3 \cdot 4 + \dots + x_n \cdot 2^{n-1}$.
In this sense, define $\overline{F}(|i\rangle) := |\overline{F}(x_1, \dots, x_n)\rangle$

An essential part of this (and many other algorithms) is the "Query":

Given 1-qubit state $|b\rangle$ and $\overline{F}(x_1, \dots, x_n) \in \{0, 1\}$, define

$$Q_{\overline{F}}(|i\rangle \otimes |b\rangle) = |i\rangle \otimes (|b\rangle \oplus \overline{F}(|i\rangle))$$

Particularly for $|b\rangle = |-\rangle := \frac{1}{\sqrt{2}}(|0\rangle - |1\rangle)$,

we have

$$\begin{aligned} Q_{\overline{F}}(|i\rangle \otimes |-\rangle) &= |i\rangle \otimes \frac{1}{\sqrt{2}} \left[\overline{F}(|i\rangle) - |1\rangle \oplus \overline{F}(|i\rangle) \right] \\ &= \begin{cases} |i\rangle \otimes \frac{1}{\sqrt{2}} [|0\rangle - |1\rangle] & \overline{F}(|i\rangle) = |0\rangle \\ |i\rangle \otimes \frac{1}{\sqrt{2}} [|1\rangle - |0\rangle] & \overline{F}(|i\rangle) = |1\rangle \end{cases} \\ &= \overline{F}(|i\rangle) \otimes |-\rangle \\ &= (-1)^{\overline{F}(|i\rangle)} |i\rangle \otimes |-\rangle \end{aligned}$$

→ Note that we identify $\overline{F}(|i\rangle)$ with $\overline{F}(x_1, \dots, x_n)$ instead of $|\overline{F}(x_1, \dots, x_n)\rangle$, i.e., the 1st equality only makes sense for basis element $|i\rangle$

Note that Q_F is unitary since

$$|b\rangle \mapsto (|b\rangle \oplus F(|a\rangle))$$

is permutation for any value of $F(a)$

We make extensive use of the Hadamard gate H given by $\frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix}$

We already know:

$$H|0\rangle = \frac{1}{\sqrt{2}}(|0\rangle + |1\rangle)$$

$$H|1\rangle = \frac{1}{\sqrt{2}}(|0\rangle - |1\rangle)$$

Lemma 1: There holds

$$\begin{aligned} H^{\otimes n} |0^n\rangle &:= \bigotimes_{i=1}^n H|0\rangle = \bigotimes_{i=1}^n \frac{1}{\sqrt{2}} (|0\rangle + |1\rangle) \\ &= \sum_{\substack{(x_1, \dots, x_n) \\ \in \{0,1\}^n}} \frac{1}{\sqrt{2^n}} |x_1 \dots x_n\rangle = \frac{1}{\sqrt{N}} \sum_{i=0}^{2^n-1} |i\rangle \end{aligned}$$

as well as for $i = x_1 + 2x_2 + \dots + 2^{n-1}x_n$

$$H^{\otimes n} |i\rangle := \bigotimes_{i=1}^n H|x_i\rangle = \bigotimes_{i=1}^n \frac{1}{\sqrt{2}} (|0\rangle + (-1)^{x_i} |1\rangle)$$

$$= \sum_{(y_1, \dots, y_n) \in \{0,1\}^n} \frac{1}{\sqrt{N}} (-1)^{x \cdot y} |y_1 \dots y_n\rangle$$

$$= \frac{1}{\sqrt{N}} \sum_{j=0}^{2^n-1} (-1)^{i \cdot j} |j\rangle, \quad \text{where}$$

$$i \cdot j := x \cdot y := \sum_{k=1}^n x_k y_k.$$

The Algorithm reads:

1) Initial state $|0^n\rangle := \underbrace{|0\rangle \otimes \dots \otimes |0\rangle}_{n\text{-times}}$

2) Apply \boxed{H} to each qbit to obtain uniform state

$$\frac{1}{\sqrt{N}} \sum_{i=0}^{2^n-1} |i\rangle$$

3) Tensorize result with $|-\rangle \Rightarrow \frac{1}{\sqrt{N}} \sum |i\rangle |-\rangle$

3) Apply query $\boxed{Q_f}$ with $|b\rangle = |-\rangle$

to obtain

$$\frac{1}{\sqrt{N}} \sum_{i=0}^{2^n-1} (-1)^{f(i)} |i\rangle |-\rangle$$

4) Ignore last qbit $|-\rangle$ and apply \boxed{H} to remaining state to obtain (Lemma 1)

$$\frac{1}{N} \sum_{i=0}^{2^n-1} \sum_{j=0}^{2^n-1} (-1)^{i \cdot j} (-1)^{f(i)} |j\rangle$$

5) The amplitude of the $j=0 \Leftrightarrow |0^n\rangle$ -state is

$$\frac{1}{N} \sum_{i=0}^{2^n-1} \underbrace{(-1)^{i \cdot 0}}_{=1} (-1)^{f(i)}$$

• If \bar{F} is balanced \Rightarrow

$$\frac{1}{N} \sum_{i=0}^{2^N-1} (-1)^{\bar{F}(i)} = 0$$

• If $\bar{F} = 1 \Rightarrow \frac{1}{N} \sum_{i=0}^{2^N-1} (-1)^{\bar{F}(i)} = -1$

• If $\bar{F} = 0 \Rightarrow \frac{1}{N} \sum_{i=0}^{2^N-1} (-1)^{\bar{F}(i)} = 1$

Complexity of Quantum Alg:

$O(1)$ steps, $O(\log_2(N))$ quantum gates

$O(\log(N))$ qbits \Rightarrow "exponential speedup"

Example for \bar{F} :

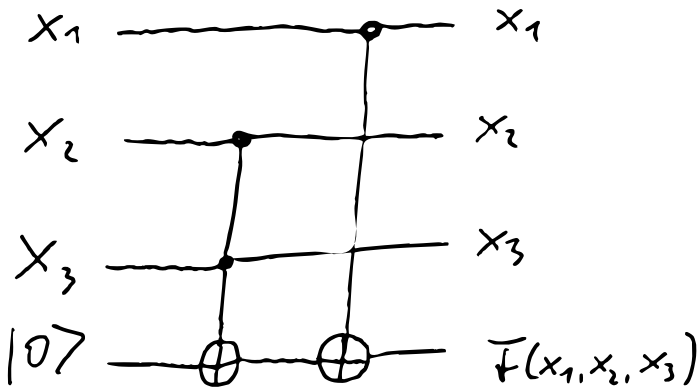
• $\bar{F}(x_1, x_2, x_3) = \text{mod}(x_1 + x_2 x_3, 2)$

$\Rightarrow \bar{F}$ is balanced since

$$\bar{F}(0, x_2, x_3) = 1 - \bar{F}(1, x_2, x_3)$$

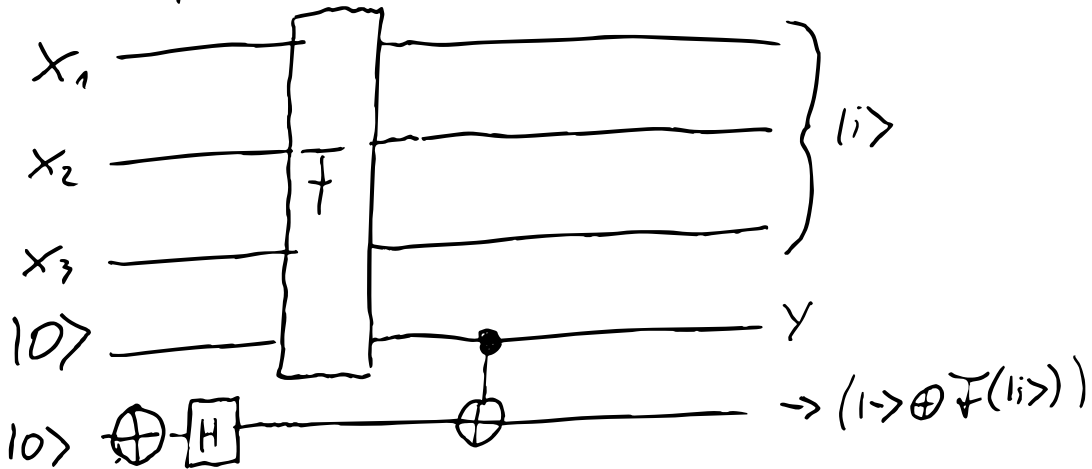
For the algorithm to work, \bar{F} needs to be a quantum circuit, i.e.

$$\bar{F}(x_1, x_2, x_3) = x_1 \text{ XOR } (x_2 \text{ AND } x_3)$$

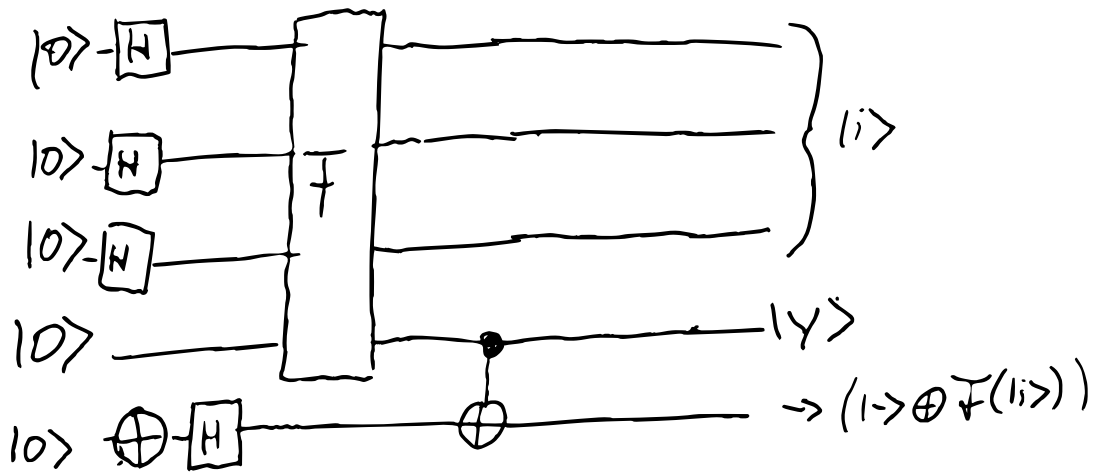


(ancilla qubit

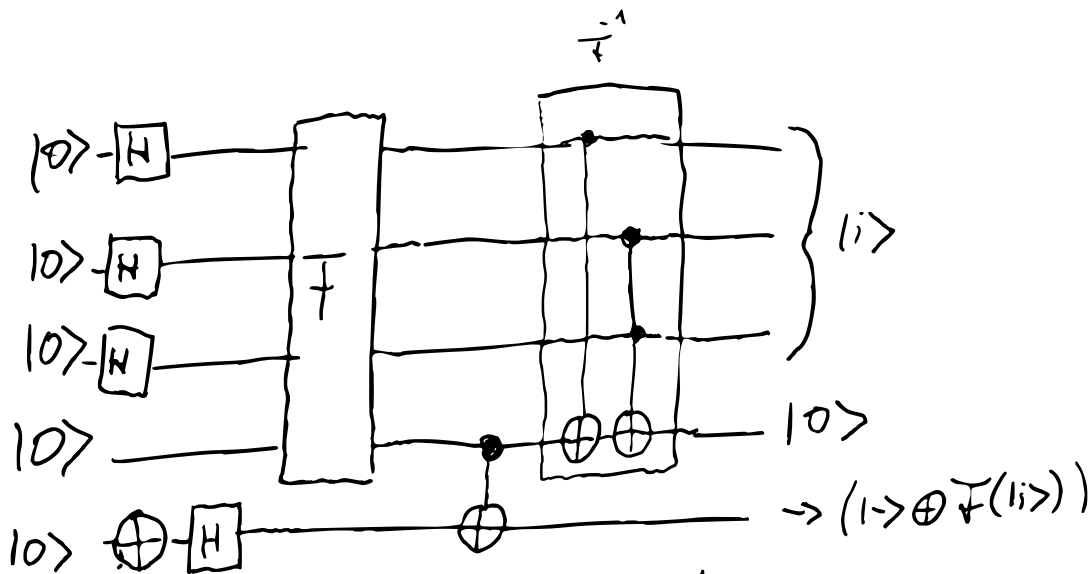
Next, implement Query $Q_{\bar{F}}(|i\rangle) = |i\rangle \oplus (|i\rangle \oplus \bar{F}(|i\rangle))$



Apply Steps 1-3 of ALP



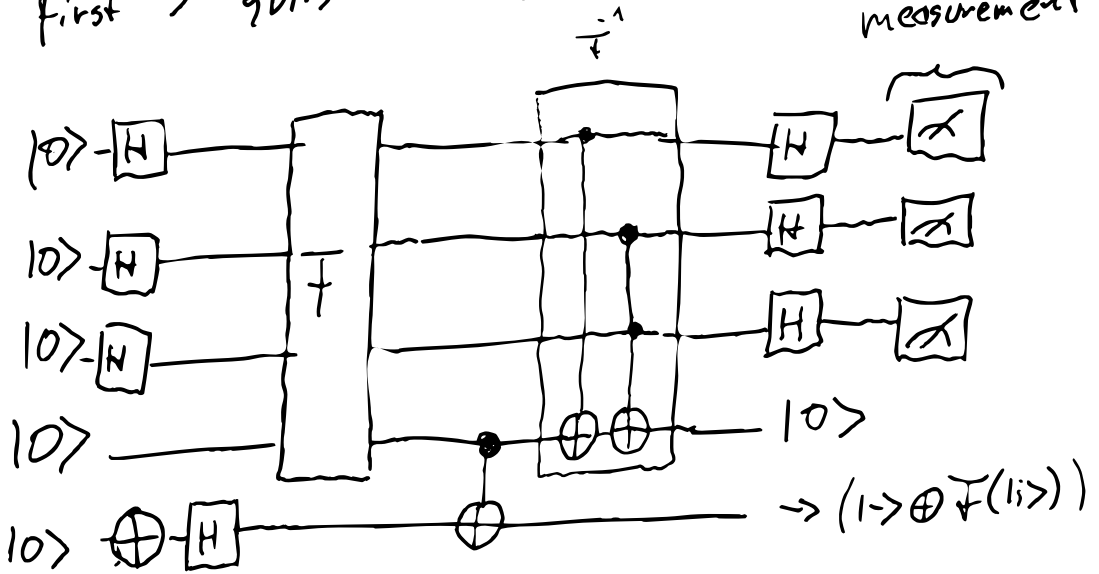
Problem: $|y\rangle$ and $|l_i\rangle$ are entangled
 Solution: reverse F such that $|y\rangle = |10\rangle$



Now, we have the desired state

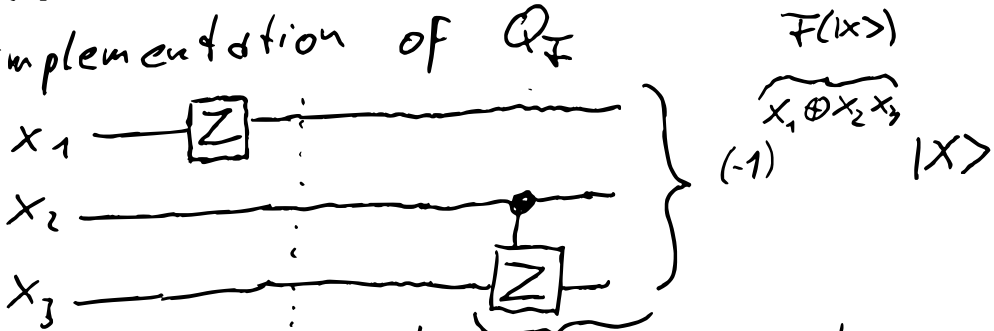
$$|l_i\rangle \otimes (|1\rangle \oplus F(|l_i\rangle)) \otimes |10\rangle$$

Apply step 4 of AIP and measure the first 3 qubits



The measurement must contain the projection onto $|000\rangle \oplus \text{span}\{|0\rangle, |1\rangle\}$

Note that this is not the most efficient implementation of Q_F



at this point we have $(-1)^{x_1} |x\rangle$

Controlled Z-gate:
Z-gate only active if $x_2=1$

$$Z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

$$\Rightarrow (-1)^{x_2 x_3} |x\rangle$$

Remarks on Deutsch-Jozsa:

• First algorithm with exponential speedup over classical computer

• However, consider the following randomized classical algorithm $R(F)$:

1) generate two random $x, y \in \{0, 1\}^n$

2) Evaluate $F(x), F(y)$

3) Output: $R(F) = 1$ if $F(x) = F(y) = 1$

$R(F) = 0$ if $F(x) = F(y) = 0$

$R(F)$ "balanced" if $F(x) \neq F(y)$

• The algorithm is correct if F is constant and answers correctly with prob. $\frac{1}{2}$ if F is balanced.

• Apply the alg k -times with i.i.d random samples to get $(R_1(F), \dots, R_k(F))$. Return

1 if $R_1(F) = \dots = R_k(F) = 1$

0 if $R_1(F) = \dots = R_k(F) = 0$

"balanced" else

\Rightarrow Probability of error 2^{-k} , cost $O(k)$

Remark: Why do you have to "uncompute F " in order to remove entanglement?

After step 3, we have

$$\frac{1}{\sqrt{N}} \sum_{i=0}^{2^n-1} (-1)^{F(i)} |i\rangle$$

but in implementation, we actually have

$$\frac{1}{\sqrt{N}} \sum_{i=0}^{2^n-1} (-1)^{F(i)} |i\rangle \otimes |0_i\rangle \otimes |-\rangle$$

For some ancilla qbit $|0_i\rangle$. We may ignore last qbit since not entangled.

Step 4 applies $H^{\otimes n}$ to obtain

$$\frac{1}{N} \sum_{i=0}^{N-1} \sum_{j=0}^{N-1} (-1)^{F(i)} (-1)^{ij} |j\rangle \otimes |0_i\rangle$$

\Rightarrow Amplitude of $|j\rangle = |0\rangle \Rightarrow$

$$\frac{1}{N} \sum_{i=0}^{N-1} (-1)^{F(i)} |0\rangle \otimes |0_i\rangle$$

Might be non-zero even for balanced F

Simon's algorithm

First quantum alg. with exp speedup over any (randomized) classical alg.

Given $i, j \in \{0, 1\}^n$, define $i \oplus j := (i_1 \oplus j_1, \dots, i_n \oplus j_n)$

Simon's problem: Let $F: \{0, 1\}^n \rightarrow \{0, 1\}^n$ s.t. there exists $s \in \{0, 1\}^n$

$$F(x) = F(y) \iff x = y \text{ or } x \oplus s = y$$

Goal: Find s . $\forall x, y \in \{0, 1\}^n$

Quantum Algorithm

1) Start with $|0^n\rangle|0^n\rangle$ and apply H to the first n qubits to obtain

$$\frac{1}{\sqrt{N}} \sum_{i=0}^{2^n-1} |i\rangle|0^n\rangle$$

2) Apply the query $|i\rangle \otimes |b^n\rangle \mapsto |i\rangle \otimes (|b^n\rangle \oplus F(i))$ with $|b^n\rangle = |0^n\rangle$ to obtain

$$\frac{1}{\sqrt{N}} \sum_{i=0}^{2^n-1} |i\rangle |F(i)\rangle$$

3) Measure the second n qubits in computational basis, i.e.

$$\text{span} \{ |i\rangle, i=0, \dots, 2^{n-1} \} \otimes \text{span} \{ |j\rangle \}$$

$j=0, \dots, 2^{n-1}$

to

obtain some output

$$|a\rangle \otimes |j\rangle = \frac{P_j \left(\frac{1}{\sqrt{N}} \sum_i |i\rangle F(|i\rangle) \right)}{\|P_j(\dots)\|}$$

Note that

$$|i\rangle F(|i\rangle) \perp \text{ran } P_j \text{ for } F(|i\rangle) \neq |j\rangle$$

by assumption, there exist exactly two inputs $|i_0\rangle$ and $|i_1\rangle = |i_0 \oplus s\rangle$ with $F(|i_k\rangle) = |j\rangle$ $k=0,1$.

$$\text{Hence } |a\rangle = \frac{1}{\sqrt{2}} (|i_0\rangle + |i_0 \oplus s\rangle)$$

4) Ignore second n qubits and apply \boxed{H} to the first n to obtain

$$\frac{1}{\sqrt{2N}} \sum_{j=0}^{2^n-1} \left[(-1)^{i \cdot j} + (-1)^{(i \oplus s) \cdot j} \right] |j\rangle$$

J) Measure in computational basis

• $|j\rangle$ has non-zero amplitude if

$$(i_0 \oplus s) \cdot j \equiv i_0 \cdot j \pmod{2} \Leftrightarrow s \cdot j \equiv 0 \pmod{2}$$

\Rightarrow We obtain random element of the set

$$\{j \mid s \cdot j \equiv 0 \pmod{2}\}$$

6) Repeat the procedure to obtain
 $n-1$ linearly independent elements $j^{(1)}, \dots, j^{(n-1)}$ with
 $j^{(i)} \cdot s = 0 \pmod{2}$

or

$$\begin{pmatrix} j^{(1)} \\ \vdots \\ j^{(n-1)} \end{pmatrix} s = 0 \pmod{2}$$

7) Solve linear system mod 2 on
classical computer in $O(n^3)$

Note that $\# \text{span} \{ j^{(1)}, \dots, j^{(k)} \} \leq 2^k$

Hence, with probability $\frac{2^n - 2^k}{2^n} = 1 - 2^{k-n} \geq \frac{1}{2}$

for $k \leq n-1$, we find

a linearly independent vector j_{k+1}

Conclusion: number of qubits $O(n)$,
number of gates $O(n)$, number of iterations
 $O(n)$ with high probability + $O(n^3)$

Classical algorithms for Simon's problem:

Lemma 2 Any (randomized) classical algorithm with less than $\delta 2^{\frac{n}{2}}$ queries to F fails with probability $\geq e^{-\frac{\delta}{\delta^2}}$. (deterministic alg. fail certainly if less than $2^{\frac{n}{2}}$ queries).

Proof Every alg generates a sequence of queries x_1, \dots, x_n with $F(x_1), \dots, F(x_n)$

If all y_i are distinct, the alg can't distinguish between $s=0$ and $s \neq 0$.

Assume all y_1, \dots, y_n are distinct. Then the alg chooses x_{n+1} based on some prob. measure μ on $\{0,1\}^n$ [in the det. case, μ is a delta distribution].

$$\sum_{k=1}^c \sum_{s \in \{0,1\}^n} \sum_{i=1}^k \mu(x_i \oplus s) = \sum_{k=1}^c \sum_{i=1}^k \underbrace{\sum_s \mu(x_i \oplus s)}_{=1} = \frac{c(c+1)}{2}$$

$$\Rightarrow \exists s \in \{0,1\}^n : \underbrace{\sum_{k=1}^c \sum_{i=1}^k \mu(x_i \oplus s)}_{p_i \leq \frac{1}{2}} \leq \frac{c(c+1)}{2^{n+1}} \leq \frac{1}{2} \quad \text{for } c(n) \leq 2^n$$

Hence,

$$\mathbb{P}(y_1, \dots, y_{n+1} \text{ distinct}) = \frac{\mathbb{P}(y_{n+1} \notin \{y_1, \dots, y_n\} | y_1, \dots, y_n \text{ distinct})}{\mathbb{P}(y_1, \dots, y_n \text{ distinct})} \geq 1 - p_i$$

Iterate the argument to obtain

$$P(\gamma_1 \dots \gamma_c \text{ distinct}) \geq \prod_{i=1}^c (1-p_i)$$

Taking logarithms shows

$$\log \left(\prod_{i=1}^c (1-p_i) \right) = \sum_{i=1}^c \log(1-p_i)$$

$$(p_i \leq \frac{1}{2}) \rightarrow \geq -\frac{5}{4} \sum_{i=1}^c p_i \geq -\frac{5}{4} \frac{c(c+1)}{2^{u+1}}$$

$$\Rightarrow P(\gamma_1 \dots \gamma_c \text{ distinct}) \geq e^{-\frac{5}{4} \frac{c(c+1)}{2^{u+1}}}$$

Choosing $c+1 = \delta 2^{\frac{u}{2}}$ satisfies $c(c+1) \leq 2^u$

and $P(\gamma_1, \dots, \gamma_c \text{ distinct}) \geq e^{-\frac{5}{4} \delta^2}$

□

$$\begin{aligned} \log(1-p) &= 0 + \frac{1}{1}(-p) - \frac{1}{2} p^2 \quad \text{for some } 1-p \leq \xi \leq p \\ &\geq -p - \frac{p^2}{2} = -p \left(1 + \frac{p}{2}\right) \\ &\geq -\frac{5}{4} p \end{aligned}$$

The quantum Fourier Transform

Discrete FT: for x_0, \dots, x_{N-1} define

$$\hat{X}_k := \frac{1}{\sqrt{N}} \sum_{j=0}^{N-1} x_j e^{-\frac{2\pi i}{N} k j} \quad \left| \quad \begin{array}{l} \text{inverse} \\ X_k = \frac{1}{\sqrt{N}} \sum_{j=0}^{N-1} x_j e^{\frac{2\pi i}{N} k j} \end{array} \right.$$

in matrix notation

$$\hat{X} = F_N X \quad \text{with} \quad F_N \in \mathbb{C}^{N \times N}$$

$$(F_N)_{kj} := \frac{1}{\sqrt{N}} e^{-\frac{2\pi i}{N} k j}$$

F_N is unitary matrix ✓

Fast FT: Assume $N = 2^n$

$$\hat{X}_k = \frac{1}{\sqrt{2}} \left(\frac{1}{\sqrt{N/2}} \sum_{\substack{j=0 \\ j \text{ even}}}^{N/2-1} x_j e^{-\frac{2\pi i}{N/2} k \frac{j}{2}} + e^{-\frac{2\pi i}{N} k} \frac{1}{\sqrt{N/2}} \sum_{\substack{j=0 \\ j \text{ odd}}}^{N/2-1} x_j e^{-\frac{2\pi i}{N/2} k \frac{j-1}{2}} \right)$$

⇒ split FT into 2 FTs of half size
leads to $O(N \log N)$ complexity.

The QFT maps a state

$$|x\rangle = \sum_{j=0}^{N-1} \hat{x}_j |j\rangle \quad \text{to} \quad |y\rangle = \sum_{j=0}^{N-1} x_j |j\rangle$$

where x_j is given by the classical IFT

If $|x\rangle = |j\rangle$, then $x_j = \delta_{jj}$ and hence

$$\hat{x}_k = \frac{1}{\sqrt{N}} e^{\frac{2\pi i}{N} k j_0}$$

$$\Rightarrow |j\rangle \xrightarrow{\text{QFT}} \frac{1}{\sqrt{N}} \sum_{k=0}^{N-1} e^{\frac{2\pi i}{N} k j} |k\rangle$$

$$= F_N |j\rangle$$

This is a convention.
Everything works with FT

To implement F_N efficiently, we observe

$$\frac{k}{N} = \frac{k}{2^n} = \sum_{e=1}^n k_e 2^{-e} \quad \text{if } k = k_{n-1} 2^{n-1} + k_{n-2} 2^{n-2} + \dots + k_0$$

$$F_N |j\rangle = \frac{1}{\sqrt{2^n}} \sum_{k \in \{0,1\}^n} e^{2\pi i j \sum_{e=1}^n k_e 2^{-e}} |k_1 \dots k_n\rangle$$

$$= \bigotimes_{e=1}^n e^{-2\pi i j k_e 2^{-e}} |k_e\rangle$$

$$= \bigotimes_{e=1}^n (|0\rangle + e^{2\pi i j / 2^e} |1\rangle) \frac{1}{\sqrt{2}}$$

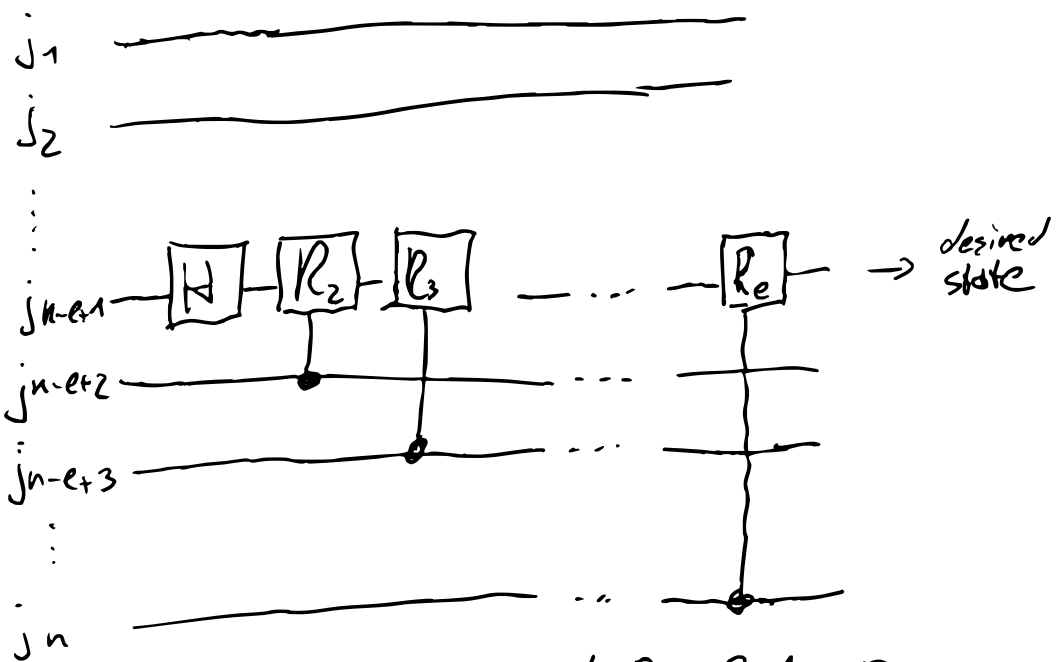
Furthermore note that

$$e^{2\pi i j 2^e} = e^{2\pi i \underbrace{\sum_{m=1}^{n-e} j_m 2^{n-m-e}}_{=1} \underbrace{2^e}_{\in \mathbb{N}}} = e^{2\pi i \sum_{m=n-e+1}^n j_m 2^{n-m-e}}$$

\Rightarrow the first $n-e$ significant bits of j do not matter.

To implement $\frac{1}{\sqrt{2}}(|0\rangle + e^{2\pi i \sum_{m=n-e+1}^n j_m 2^{n-m-e}} |1\rangle)$ we use the R_3 gate or P in IBM q-Composer

given by the matrix $\begin{pmatrix} 1 & 0 \\ 0 & e^{2\pi i \frac{2^e}{2^2}} \end{pmatrix}$



need to repeat that for $e=1, \dots, l_2$

Cost of QFT:

We need n qubits and $O(n)$ gates per qubit

$\Rightarrow O(n^2)$ gates

\Rightarrow exponential speedup over FFT with $O(n 2^n)$ operations.

Remark Strictly speaking, QFT does something different than FT. The state

$$\text{QFT}(|x\rangle) = \sum_{j=0}^{N-1} x_j |j\rangle$$

can only be accessed via measurement and hence will collapse to some $|j'\rangle$ with a certain probability. We will see that this is still very useful.

Remark R_s gates don't do very much for large s . One can show that $O(n \log n)$ gates suffice if one accepts a small error probability

Remark reversing the order of the gates and using the adjoint gates gives an efficient implementation of the inverse QFT F_N^{-1}

Application Phase estimation

Suppose we have unitary operator $U: V \rightarrow V$
with eigenvector ψ
 $\dim V = 2^n$

$$\Rightarrow U\psi = e^{2\pi i \phi} \psi \quad \text{for some } \phi \in [0, 1)$$

$$(|U\psi|^2 = \lambda^2 |\psi|^2 = |\psi|^2 \Rightarrow |\lambda| = 1)$$

Assume that $\phi = \sum_{j=1}^n \phi_j 2^{-j}$ can be written with n bits

Since U is operator on 2^n -dim Hilbert space classical computation of $U\psi \cdot \psi$ costs at least $O(2^n)$

Quantum algorithm

- 1) Start with $|0^n\rangle |\psi\rangle$
- 2) Apply $H^{\otimes n}$ to the first n qubits to obtain $\frac{1}{\sqrt{N}} \sum_{j=0}^{N-1} |j\rangle |\psi\rangle$ (Applying QFT would do the same)
- 3) Apply $|j\rangle |\psi\rangle \mapsto |j\rangle U^j |\psi\rangle$ to obtain $\frac{1}{\sqrt{N}} \sum_{j=0}^{N-1} e^{2\pi i j \phi} |j\rangle |\psi\rangle$

Note that the first n qubits satisfy

$$\frac{1}{\sqrt{N}} \sum_{j=0}^{N-1} e^{2\pi i j \phi} |j\rangle = \frac{1}{\sqrt{N}} \sum_{j=0}^{N-1} e^{2\pi i j \frac{N\phi}{N}} |j\rangle = \mathbb{F}_N(|N\phi\rangle)$$

→ This might be a bit confusing since suddenly $\phi \in \mathbb{R}$ is interpreted as an element of the vector space in which ψ is contained

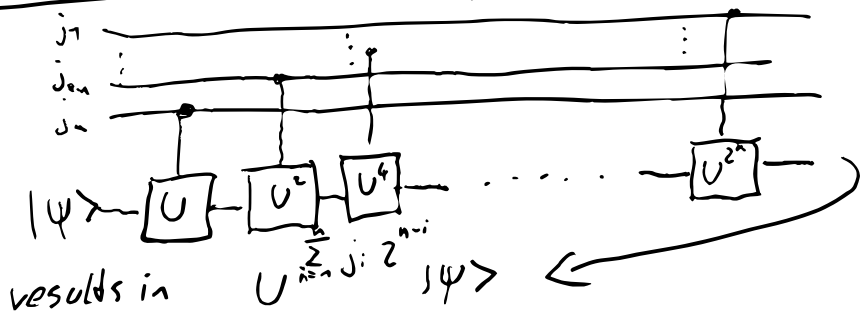
But $N\phi = \sum_{j=1}^n \phi_j 2^{n-j}$ is a basis element and hence this makes sense

4) Apply IQFT to the first n qubits to obtain

$$\mathbb{F}_N^{-1} \left(\frac{1}{\sqrt{N}} \sum_{j=0}^{N-1} e^{2\pi i j \phi} |j\rangle \right) = |N\phi\rangle$$

5) Measure in computational basis to obtain $N\phi$ and hence ϕ

Remark How to implement step 3?



But we will need to assume that U^{2^k} can be implemented effectively (This might not be true in general, but is in the applications below)

Note that the input doesn't need to be a single eigenvector. Let $|\psi\rangle, |\psi'\rangle$ denote two EVs with phase ϕ, ϕ' . Linearity implies that input $\frac{1}{\sqrt{2}}(|\psi\rangle + |\psi'\rangle)$ produces output $\frac{1}{\sqrt{2}}(|N\phi\rangle + |N\phi'\rangle)$

A measurement will produce ϕ, ϕ' with equal probability

if ϕ requires more than n bits, the final state is a small perturbation δ of $T_n(|R\rangle)$ with

$$|\delta| = \left| \frac{1}{\sqrt{N}} \sum_{j=0}^{N-1} \underbrace{\begin{bmatrix} 2\alpha^{ij\phi} & 2\alpha^{ij\phi'} \\ e & -e \end{bmatrix}}_{O(2^{-n})} |j\rangle \right| \left| \frac{1}{\sqrt{N}} \sum e^{2\alpha^{ij\phi}} |j\rangle \right| \approx O(2^{-n})$$

Show's integer factorization

Factoring problem: given $n \in \mathbb{N} \setminus \{\text{primes}\}$
Find $1 < k < n$ with $\frac{n}{k} \in \mathbb{N}$.

- $n \in \mathbb{N}$ can be defined using $\log_2(n)$ bits \Rightarrow Poly. complexity of $\log_2(n)$ means $O(\log_2(n)^P)$ operations
- There exist efficient ^{classical} algorithms to check whether n is prime (see, e.g. AKS-test or BPSW-test) or power of prime
- Security of many crypto-systems is based on the fact that integer factorization is hard
- Best known classical algorithm is the general number field sieve with runtime $O(e^{(\log_2 n)^{\frac{1}{3}} (\log_2 \log_2 n)^{\frac{2}{3}}})$

Note that the runtime is a conjecture!

- It is not known that classical algorithms can not be faster. Latest paper which (falsely) claimed an $O(\log n^3)$ -alg for factorization was by Schnorr (Pöschel) in 2021

Reduction to period finding

Want to factor $N \in \mathbb{N}$.

1) Choose random $x \in \{2, \dots, N-1\}$ with $\gcd(x, N) = 1$

$\Rightarrow x \in \left(\frac{\mathbb{Z}}{N\mathbb{Z}} \right)^\times$ multiplicative group mod n

(Lemma 3)

$\Rightarrow x$ has period r with $x^r \bmod N = 1$

\Rightarrow with prob. $\frac{1}{2}$ r is even and

$$x^{\frac{r}{2}} \pm 1 \not\equiv 0 \pmod{N}$$

$$\Rightarrow (x^{\frac{r}{2}} - 1)(x^{\frac{r}{2}} + 1) \equiv x^r - 1 \equiv 0 \pmod{N}$$

$\Rightarrow \gcd(x^{\frac{r}{2}} - 1, N), \gcd(x^{\frac{r}{2}} + 1, N)$
are non-trivial factors of N

Lemma 3 Every $x \in (\mathbb{Z}/N\mathbb{Z})^\times$ has period v which is minimal $v \in \mathbb{N}$ with $x^v \equiv 1 \pmod{N}$.

Proof Consider the set $S = \{1, x, x^2, \dots\} \subseteq (\mathbb{Z}/N\mathbb{Z})^\times$

Since $(\mathbb{Z}/N\mathbb{Z})^\times$ is finite, there exist $j \neq k \in \mathbb{N}$ with $x^j = x^k \pmod{N}$. W.L.O.P assume $j < k$

$\Rightarrow x^{k-j} = 1 \pmod{N}$. □

Euler totient func: $\varphi(n) := \#\{1 \leq k \leq n : \gcd(k, n) = 1\}$

$\varphi(n) = n \prod_{\substack{p|n \\ p \text{ prime}}} (1 - \frac{1}{p}) \Rightarrow \varphi(p^a) = p^a (1 - \frac{1}{p}) = p^{a-1} (p-1)$

Lemma [Chinese remainder Thm]

Let $m_1, \dots, m_n \in \mathbb{N}$ with $\gcd(m_i, m_j) = 1 \quad i \neq j$. Then

$$\left[\begin{array}{l} \text{Find } x \text{ s.t.} \\ x \equiv a_1 \pmod{m_1} \\ \vdots \\ x \equiv a_n \pmod{m_n} \end{array} \right]$$

is solvable and any two solutions are equal mod $M = m_1 m_2 \dots m_n$

Proof Define $M_i := \frac{M}{m_i} \Rightarrow \gcd(M_i, m_i) = 1$

$\Rightarrow M_i$ has mult. inverse mod m_i denoted by N_i

$\Rightarrow x = \sum_i a_i M_i N_i$ is solution since

$M_i N_i \equiv 1 \pmod{m_i}$ and $M_i N_i \equiv 0 \pmod{m_j} \quad j \neq i$

Two solutions x, x' satisfy $x - x' \equiv 0 \pmod{m_i}$

$\Rightarrow M = m_1 m_2 \dots m_n$ divides $x - x'$ (since m_i are coprime) □

Lemma 4 Let $p > 2$ prime. Let 2^d maximal power of 2 dividing $\varphi(p^k)$. Let x denote randomly chosen element in $\left(\frac{\mathbb{Z}}{p^k \mathbb{Z}}\right)^\times$.

$$\Rightarrow P(2^d \text{ divides order of } x) = \frac{1}{2}$$

Proof: $\varphi(p^k) = p^{k-1}(p-1)$ is even $\Rightarrow d \geq 1$.

Let g denote generator of $\left(\frac{\mathbb{Z}}{p^k \mathbb{Z}}\right)^\times \Rightarrow x = g^k \pmod{p^k}$ for some $k \in \{1, \dots, \varphi(p^k)\}$. Let r denote order of x .

1) k is odd: $g^{kr} \equiv 1 \pmod{p^k} \Rightarrow \varphi(p^k)$ divides kr
 since k is odd $\Rightarrow 2^d$ divides r } since $\varphi(p^k)$ is minimal with $g^{\varphi(p^k)} = 1$

2) k is even: $g^{k\varphi(p^k)/2} = 1 \pmod{p^k}$
 $\Rightarrow r$ divides $\varphi(p^k)/2$ since $x^r = g^{kr} = 1 \pmod{p^k}$
} r is minimal with

$\Rightarrow 2^d$ does not divide r

Hence, for exactly half of $x \in \left(\frac{\mathbb{Z}}{p^k \mathbb{Z}}\right)^\times$ we have $x = g^k$ with k even/odd □

Lemma 5 Let $N = p_1^{d_1} \dots p_m^{d_m}$ prime factorization of odd N . Let x be random in $\left(\frac{\mathbb{Z}}{N\mathbb{Z}}\right)^\times$ with order $r \pmod{N}$.

$$\Rightarrow \mathbb{P}(r \text{ is even and } x^{\frac{r}{2}} \equiv -1 \pmod{N}) \geq 1 - 2^{-m+1}$$

Proof Chinese remainder theorem \Rightarrow choose x random is equivalent to choose x_j rand in $\left(\frac{\mathbb{Z}}{p_j^{d_j}}\right)^\times$ with $x = x_j \pmod{p_j^{d_j}} \quad \forall j=1, \dots, m$

(This is because $x \leftrightarrow (x_1, \dots, x_m)$ is one-to-one)

Let r_j be order $x_j \pmod{p_j^{d_j}}$ and let 2^{d_j} max power of 2 dividing r_j . $\left\{ \begin{array}{l} r \text{ is odd or } x^{\frac{r}{2}} \equiv -1 \pmod{N} \Rightarrow \\ d_1 = d_2 = \dots = d_m \end{array} \right.$

We will show $d_1 = d_2 = \dots = d_m$. If this is the case, the following argument concludes the proof:

Let d_j' maximal s.t. $2^{d_j'}$ divides $\varphi(p_j^{d_j})$

$$\text{Lemma 4} \Rightarrow \mathbb{P}(d_j = d_j') = \frac{1}{2} \Rightarrow \mathbb{P}(d_j = k) \leq \frac{1}{2}$$

$\forall k \in \mathbb{N}$. Since d_j independent

$$\Rightarrow \mathbb{P}(d_1 = \dots = d_m) \leq 2^{-m+1}$$

It remains to show $d_1 = \dots = d_m$.

Case 1: r is odd. $x^r = 1 \pmod{N} \Rightarrow x^r = 1 \pmod{p_j^{d_j}}$
 $\Rightarrow r_j$ divides $r \Rightarrow r_j$ is odd $\Rightarrow d_j = 0$.

Case 2: r is even and $x^{\frac{r}{2}} = -1 \pmod{N}$
 $\Rightarrow x^{\frac{r}{2}} = -1 \pmod{p_j^{d_j}}$. If r_j would divide $\frac{r}{2}$, we would
 have $x^{\frac{r}{2}} = x^{kr_j} = -1 \pmod{p_j^{d_j}}$
 $= x_j^{kr_j} \pmod{p_j^{d_j}}$
 $= 1 \pmod{p_j^{d_j}}$ } $\Rightarrow r_j$ does not divide $\frac{r}{2}$.

but r_j divides r . Hence largest power
 of 2 dividing r must be equal to d_j . □

Note that if r is period of $x \pmod{N}$
 and r even, we have $x^{\frac{r}{2}} \neq 1 \pmod{N}$

$\Rightarrow \mathcal{P}(r \text{ even } \mid x^{\frac{r}{2}} \neq \pm 1 \pmod{N}) \geq 1 - 2^{-m+1} \geq \frac{1}{2}$
 if $m \geq 2$, i.e., if N is not a prime power.

Define $U|y\rangle := |xy \bmod N\rangle$

Note that $|j\rangle \mapsto |xj \bmod N\rangle$ is bijective

since $xj = xj' \bmod N \implies x(j-j') = 0 \bmod N$
 $\implies \gcd(x, N) \neq 1$.

$\implies U$ is permutation on basis elements \implies unitary

Lemma Let r denote the period of x in $(\mathbb{Z}/N\mathbb{Z})^\times$.

Then, $|u_s\rangle := \frac{1}{\sqrt{r}} \sum_{k=0}^{r-1} \exp\left(\frac{-2\pi i s k}{r}\right) |x^k \bmod N\rangle$

are EVs of U with eigenvalues

$$\lambda_s := \exp\left(\frac{2\pi i s}{r}\right), \quad s = 0, \dots, r-1$$

Proof

$$U|u_s\rangle = \frac{1}{\sqrt{r}} \sum_{k=0}^{r-1} \exp\left(\frac{-2\pi i s k}{r}\right) |x^{k+1} \bmod N\rangle$$

$$= \lambda_s \frac{1}{\sqrt{r}} \sum_{k=0}^{r-1} \exp\left(\frac{-2\pi i s (k+1)}{r}\right) |x^{k+1} \bmod N\rangle$$

$$\begin{array}{l} x^r = x^0 \bmod N \\ -2\pi i s \\ e = e \end{array} \implies \lambda_s |u_s\rangle$$

□

Note that there holds

$$\begin{aligned} \frac{1}{\sqrt{r}} \sum_{s=0}^{r-1} |u_s\rangle &= \frac{1}{r} \sum_{s=0}^{r-1} \sum_{k=0}^{r-1} e^{\frac{-2\pi i k s}{r}} |x^k \bmod N\rangle \\ &= \frac{1}{r} \sum_{k=0}^{r-1} \underbrace{\sum_{s=0}^{r-1} e^{\frac{-2\pi i k s}{r}}}_{\begin{cases} 0 & k \neq 0 \\ r & k = 0 \end{cases}} |x^k \bmod N\rangle \end{aligned}$$

$$= |x^0 \bmod N\rangle = |1\rangle$$

\Rightarrow Quantum phase estimation ^{for $\sqrt{\cdot}$} with input $|1\rangle$ produces the state

$$\frac{1}{\sqrt{r}} \sum_{s=0}^{r-1} |z^{\frac{s}{r}}\rangle$$

note that $r \leq N$ hence $\frac{s}{r}$ can be exactly represented with n bits

• To efficiently implement $U^{2^k} |y\rangle = |xy^{2^k} \bmod N\rangle$

we use the fact k -times

$$y^{2^k} = \left(\left(\left(y^2 \right)^2 \right)^2 \right)^2 \dots \text{Here, } |xy^{2^k} \bmod N\rangle$$

requires 1 multiplication, k squares mod N

\Rightarrow can be implemented as in classical circuits with XOR & AND gates.

Shor's period finding algorithm

Note that this is the only quantum part of Shor's algo. $N \leq 2^n$

1) Prepare $|\psi\rangle = |1\rangle = \overbrace{[0 \dots 0]_n}^n |1\rangle$ and U as before

2) Use quantum phase estimation with U and $|\psi\rangle$ to obtain the state

$$\frac{1}{\sqrt{r}} \sum_{s=0}^{r-1} |2^n \frac{s}{r}\rangle$$

3) Measurement gives a random number $\frac{s}{r}$ from the set $\{\frac{1}{r}, \dots, \frac{r-1}{r}\} \subseteq \{\frac{i}{2^n} : i=0, \dots, 2^n-1\}$

4) Write $\frac{s}{r} = \frac{s_0}{r_0}$ with $\gcd(s_0, r_0) = 1$. If s, r were coprime already, we have $r=r_0$ and found the period. Otherwise, repeat.

Remark: There are $O(\frac{r}{\phi(r)})$ numbers $1 \leq s < r$ with $\gcd(s, r) = 1$. — Euler's totient function

\Rightarrow Success prob of step 4 is $O(\frac{1}{\phi(r)})$

Since $\frac{1}{\phi(r)} \geq \frac{1}{\phi(N)} \geq \frac{1}{\phi^n}$, step 4 requires on average ϕ^n tries.

Remark

Implicitely, we used QFT to
find the frequency of

$x^k \bmod N$, i.e. the period of x

Grover's algorithm

Problem: given $F: \{0,1\}^n \rightarrow \{0,1\}$, find $x \in \{0,1\}^n$ with $F(x)=1$ or determine $F=0$.

Recall the query

$$Q_F(|i\rangle) = (-1)^{F(i)} |i\rangle$$

(recall we identify $x \in \{0,1\}^n$ with $i = 0, \dots, 2^{n-1}$)

Define the Grover diffusion operator

$$U_s := 2|s\rangle\langle s| - \mathbb{I}$$

Projection onto $|s\rangle$

where $|s\rangle$ denotes the uniform state

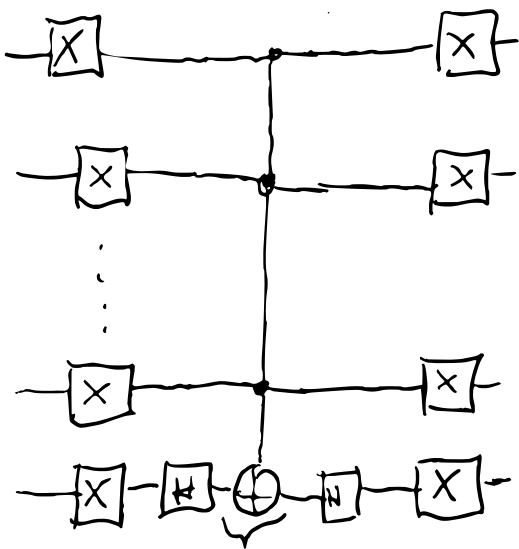
$$|s\rangle = \frac{1}{\sqrt{N}} \sum_{j=0}^{N-1} |j\rangle$$

Note that $|s\rangle = H^{\otimes n} |0^n\rangle$ and hence

$$U_s = H^{\otimes n} \left(2|0^n\rangle\langle 0^n| - \mathbb{I} \right) H^{\otimes n}$$

corresponds to the matrix

$$\begin{pmatrix} +1 & & & 0 \\ & -1 & & \\ & & \ddots & \\ 0 & & & -1 \end{pmatrix} \in \mathbb{R}^{2^n \times 2^n}$$



$$-(|210\rangle\langle 01 - I)$$

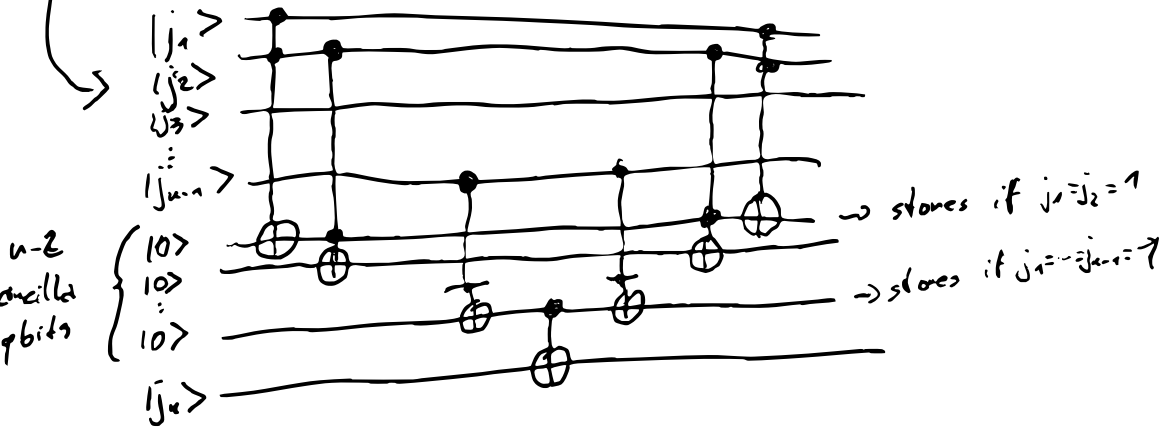
since we can only measure amplitudes, a global phase change (multiplication with $\alpha \in \mathbb{C}, |\alpha|=1$) is not noticeable alternatively, one

can apply $|Y\rangle\langle X|Z\rangle\langle X|$ to obtain

$$\begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} = \begin{pmatrix} -1 & 0 \\ 0 & -1 \end{pmatrix}$$

generalized Toffoli gate

$$\begin{pmatrix} 1 & 0 & 1 \\ 0 & 1 & 0 \end{pmatrix}$$



Algorithm

1) Apply $H^{\otimes n}$ to obtain $|S\rangle = \frac{1}{\sqrt{N}} \sum_{j=0}^{N-1} |j\rangle$

2) For $k=1, \dots, r(N)$ do

2a) Apply Q_F to current state $\frac{1}{\sqrt{N}} \sum_{j=0}^{N-1} \alpha_j |j\rangle$
to obtain $\sum_{j=0}^{N-1} (-1)^{F(j)} \alpha_j |j\rangle$

2b) Apply U_S to first n qubits

3) Measure in computational basis

Thm: With $r(N) = \frac{\arccos(\sqrt{\frac{1-\epsilon}{N}})}{2 \arcsin(\sqrt{\frac{1-\epsilon}{N}})}$, Grover's alg. finds state $|x\rangle$ with $F(x)=1$ with probability $\geq 1 - \frac{\epsilon}{N}$

Proof:

Define $|B\rangle := \frac{1}{\sqrt{N-t}} \sum_{F(j)=0} |j\rangle$, where $t = \# F^{-1}(1)$

Note that the first n qubits of $|x\rangle \mapsto Q_F(|x\rangle)$

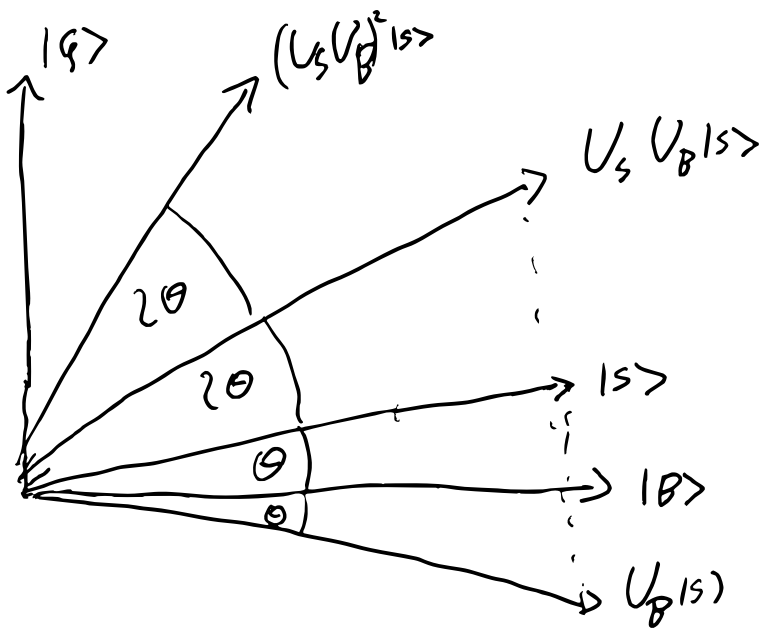
satisfy $|x\rangle \mapsto (2|B\rangle\langle B| - I)|x\rangle = U_B(|x\rangle)$

(Check for basis elements $|x\rangle = |j\rangle$)

Define $|G\rangle := \frac{1}{\sqrt{t}} \sum_{F(j)=1} |j\rangle$ and note

$|S\rangle \in \text{span}\{|G\rangle, |B\rangle\}$; $U_S, U_B: \text{span}\{|G\rangle, |B\rangle\} \rightarrow \text{span}\{|G\rangle, |B\rangle\}$

hence, the Grover iteration never leaves the plane $\text{span}\{|G\rangle, |B\rangle\}$



Each application of $U_S U_B$ rotates a state $|x\rangle$ towards $|g\rangle$ by an angle θ given by $\text{span}\{|s\rangle, |g\rangle\}$

$$\cos \theta = \frac{\langle s|B\rangle}{|s||B|} = \frac{1}{\sqrt{N}} \frac{1}{\sqrt{N+1}} \sum_{F(j)=0} \langle s|j\rangle = \frac{N+1}{\sqrt{N(N+1)}}$$

$$= \sqrt{1 - \frac{1}{N}}$$

$$\Rightarrow \sin^2 \theta = 1 - \cos^2 \theta = 1 - \left(1 - \frac{1}{N}\right) \Rightarrow \sin \theta = \sqrt{\frac{1}{N}}$$

$$\Rightarrow \theta \approx \sqrt{\frac{1}{N}}$$

For large N .

Since the angle between $|g\rangle$ and $|s\rangle$ is $\arccos\left(\frac{\sqrt{1/N}}{1}\right)$

$$\cos \alpha = \frac{1}{\sqrt{N}} \frac{1}{1} \sum_{F(j)=1} \langle s|j\rangle = \frac{\sqrt{1/N}}{1} \Rightarrow r(N) = \text{round}\left(\frac{\arccos\left(\frac{\sqrt{1/N}}{1}\right)}{\sqrt{1/N}}\right)$$

to obtain a state $|x\rangle$ with $\langle g|x\rangle \geq \cos(\theta)$

$$\Rightarrow |\langle g|x\rangle|^2 \geq \cos^2 \theta = 1 - \frac{t}{N}$$



Remark in practice, we don't know when $|x\rangle$ gets close to $|g\rangle$. But we can measure and check whether $F(x) = 1$. If not, restart the algorithm. If $t \ll N$

$$r(N) = \frac{\arccos \sqrt{\frac{t}{N}}}{\arcsin \sqrt{\frac{t}{N}}} \approx \frac{\pi}{4} \sqrt{\frac{N}{t}}$$

Remark If $t=1$, a classical algorithm requires at least N evaluations of F . Grover's algorithm only needs $O(\sqrt{N})$ iterations with $O(n)$ gates

Optimality of Prover's AIP

Lemma Any Quantum AIP based on the Query Q_F requires at least $O(\delta\sqrt{N})$ applications of Q_F to succeed with prob. $\geq \delta^2$.

Proof Any Q-AIP starts with some state $|\psi\rangle$ and applies Unitary transformations as well as Q_F . i.e. we may write the state after k applications of Q_F as

$$|\psi_k^F\rangle := \underbrace{U_k Q_F U_{k-1} Q_F \dots U_1 Q_F}_{\text{unitary}} |\psi\rangle$$

Additionally, we will consider

$$|\psi_k\rangle := U_k U_{k-1} \dots U_1 |\psi\rangle$$

Let $F_j: \{0,1\}^n \rightarrow \{0,1\}$ with $F_j(|i\rangle) = \delta_{ij}$

and define

$$D_k = \sum_{j=0}^{N-1} \|\psi_k^{F_j} - \psi_k\|^2$$

Idea: If D_k is small, the evolution of F doesn't make a big difference and it will be hard to find $F(|i\rangle) = 1$.

Step 1: Show that $D_k \leq 4k^2$ by induction.

$k=0$: $D_k = 0$ ✓

$k \rightarrow k+1$:
$$D_{k+1} = \sum_{j=0}^{N-1} \|U_{k+1} Q_{F_j} \psi_k^{F_j} - V_{k+1} \psi_k\|^2$$

$$= \sum_{j=0}^{N-1} \|Q_{F_j} \psi_k^{F_j} - \psi_k\|^2$$

$$= \sum \|Q_{F_j} (\psi_k^{F_j} - \psi_k) + (Q_{F_j} - I) \psi_k\|^2$$

Note that $Q_{F_j} = I - 2|j\rangle\langle j| \Rightarrow (Q_{F_j} - I)|\psi_k\rangle = -2|j\rangle\langle j|\psi_k\rangle$

$$= \underline{-2|j\rangle\langle j|\psi_k\rangle}$$

\Rightarrow

$$D_{k+1} \leq \sum_{j=0}^{N-1} \|\psi_k^{F_j} - \psi_k\|^2 + 4\|\psi_k^{F_j} - \psi_k\| |\langle j|\psi_k\rangle| + 4|\langle j|\psi_k\rangle|^2$$

Cauchy - Schwarz shows

$$D_{k+1} \leq D_k + 4 \left(\sum_j \|\psi_k^{\bar{F}_j} - \psi_k\|^2 \right)^{\frac{1}{2}} \underbrace{\left(\sum_j |\langle j | \psi_k \rangle|^2 \right)^{\frac{1}{2}}}_{=1}$$

$$+ 4 \underbrace{\langle \psi_k | \psi_k \rangle^2}_{=1}$$

$$\leq D_k + 4\sqrt{D_k} + 4$$

Induction Hyp. $\rightarrow \leq 4k^2 + 8k + 4 = \underline{4(k+1)^2}$ ✓

This concludes the induction and shows

$$\underline{D_k \leq 4k^2}$$

Step 2: Assume $|\langle j | \psi_k^{\bar{F}_j} \rangle|^2 \geq \delta > 0 \forall j=0, \dots, k-1$, i.e. the alg. works with high prob. for each input.

Replacing $|j\rangle$ with $e^{i\theta_j} |j\rangle$ does not change success prob. \Rightarrow we may assume $|\langle j | \psi_k^{\bar{F}_j} \rangle| = |\langle j | \psi_k^{\bar{F}_j} \rangle|$.

$$\Rightarrow \|\psi_k^{\bar{F}_j} - j\|^2 = 2 - 2|\langle j | \psi_k^{\bar{F}_j} \rangle| \leq \underline{2(1-\delta)}$$

Define $E_k := \sum_{j=0}^{k-1} \|\psi_k^{\bar{F}_j} - j\|^2 \Rightarrow E_k \leq 2N(1-\delta)$

$$F_k := \sum_{j=0}^{k-1} \|j - \psi_k\|^2$$

$$\begin{aligned}
\Rightarrow \\
D_k &= \sum_{j=0}^{N-1} \|(\psi_k^{\dagger j} - j) + (j - \psi_k)\|^2 \\
&\geq \sum_j \|\psi_k^{\dagger j} - j\|^2 - 2 \|\psi_k^{\dagger j} - j\| \|j - \psi_k\| \\
&\quad - \|j - \psi_k\|^2 \\
&= E_k + \bar{F}_k - 2 \left(\sum_j \|\psi_k^{\dagger j} - j\| \right)^{\frac{1}{2}} \left(\sum_j \|j - \psi_k\|^2 \right)^{\frac{1}{2}} \\
&= E_k + \bar{F}_k - 2 \sqrt{E_k \bar{F}_k} = \underline{\underline{(\sqrt{E_k} - \sqrt{\bar{F}_k})^2}}
\end{aligned}$$

Note that any state $|\phi\rangle$ satisfies

$$\begin{aligned}
\sum_{j=0}^{N-1} \|\phi - j\|^2 &= \sum_j \|\phi\|^2 - 2 \langle \phi | j \rangle + \|j\|^2 \\
&\geq 2N - 2 \sqrt{\sum_j 1} \underbrace{\sqrt{\sum_j |\langle \phi | j \rangle|^2}}_{=1} \\
&= \underline{\underline{2(N - \sqrt{N})}}
\end{aligned}$$

This implies $\bar{F}_k \geq 2(N - \sqrt{N})$ and hence

$$\begin{aligned}
\bar{F}_k \geq E_k \text{ for sufficiently large } N \text{ and} \\
\sqrt{\bar{F}_k} - \sqrt{E_k} \geq \sqrt{2(N - \sqrt{N})} - \sqrt{2(1 - \delta)N} = \frac{2\delta N - 2\sqrt{N}}{\sqrt{2(N - \sqrt{N})} + \sqrt{2(1 - \delta)N}}
\end{aligned}$$

$$\geq \frac{\epsilon \delta N - \epsilon \sqrt{N}}{\epsilon \sqrt{2N}} = \frac{\delta}{\sqrt{2}} \sqrt{N} - \frac{1}{\sqrt{2}}$$

$$\Rightarrow D_k \geq (\sqrt{F_k} - \sqrt{E_k})^2 \approx \delta^2 N$$

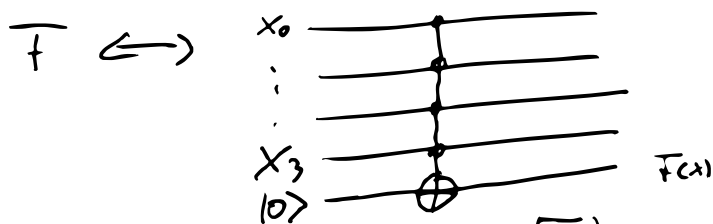
$$\text{Since } D_k \leq 4k^2 \Rightarrow k \approx \delta \sqrt{N}$$

□

Example in Quantum Composen

Link on webpage, $N=16, u=4$

$$\begin{aligned} F(x) = F(x_0, \dots, x_3) &= x_0 \text{ AND } x_1 \text{ AND } x_2 \text{ AND } x_3 \\ &= \delta_{x=(1,1,1,1)} \Rightarrow t=1 \end{aligned}$$



There holds $\frac{\arccos(\sqrt{\frac{1}{16}})}{2 \arcsin(\sqrt{\frac{1}{16}})} \approx 2.6083$

\Rightarrow optimal $r(N) = 3$.

$$\text{Probability of success} \geq 1 - \frac{1}{N} = \frac{15}{16} \approx 0.9375$$

Numerical Quadrature:

Problem: Given $f: \{0,1\}^n \rightarrow [-1,1]$, compute $\frac{1}{N} \sum_{i=0}^{N-1} f(i)$

Quantum Super Sampling

Assumption: Oracle $Q_f: |0\rangle \otimes |i\rangle \mapsto \left[\sqrt{1-f(i)} |0\rangle + \sqrt{f(i)} |1\rangle \right] \otimes |i\rangle$

1) Start with $|0^p\rangle \otimes |0\rangle \otimes |0^n\rangle$

2) Apply $H^{\otimes p} \otimes I \otimes H^{\otimes n} \rightarrow$

$$\frac{1}{\sqrt{PN}} \sum_{i=0}^{N-1} \sum_{m=0}^{P-1} |m\rangle \otimes |0\rangle \otimes |i\rangle$$

3) Apply $Q_f \rightarrow$

$$\frac{1}{\sqrt{PN}} \sum_i \sum_m |m\rangle \otimes \left[\sqrt{1-f(i)} |0\rangle + \sqrt{f(i)} |1\rangle \right] \otimes |i\rangle$$

4) Define $|G\rangle := \frac{1}{\sqrt{\sum_i f(i)}} \sum_i \sqrt{f(i)} |1\rangle \otimes |i\rangle$

$$|B\rangle := \frac{1}{\sqrt{\sum_i (1-f(i))}} \sum_i \sqrt{1-f(i)} |0\rangle \otimes |i\rangle$$

$$\bar{f} := \frac{1}{N} \sum_i f(i)$$

Note

$$U_s := (|210\rangle\langle 01| - I) \otimes I$$

$$U_s |B\rangle = |B\rangle$$

$$U_s |G\rangle = -|G\rangle$$

$$U_\psi = (|21\psi\rangle\langle\psi| - I)$$

$$\text{with } |\psi\rangle := \sqrt{1-p} |B\rangle + \sqrt{p} |G\rangle$$

Steps 1-3 provide quantum circuit to implement

the map $U: |0\rangle \otimes |0^n\rangle \mapsto |\psi\rangle$

$$\Rightarrow U_\psi = \bar{U}^{-1} (|210^{n+1}\rangle\langle 0^{n+1}| - I) U$$

as in Grover's Algorithm.

5) state reads

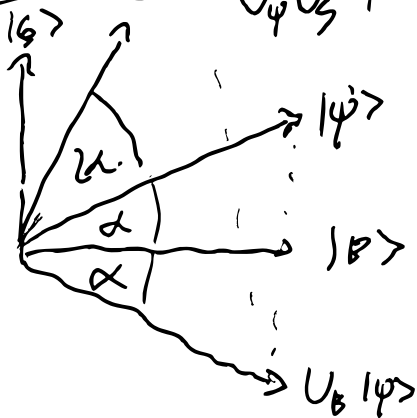
$$\frac{1}{\sqrt{P}} \sum_{m=0}^{P-1} |m\rangle \otimes [\sqrt{1-p} |B\rangle + \sqrt{p} |G\rangle]$$

for $\bar{p} := \frac{1}{N} \sum f(i)$

6) Apply prover iteration $U_\psi U_S$ m -times to second $n+1$ qbits to obtain

$$\frac{1}{\sqrt{P}} \sum_{m=0}^{P-1} |m\rangle \otimes (U_\psi U_S)^m [\sqrt{1-p} |B\rangle + \sqrt{p} |G\rangle]$$

Explanation:



Let $\sin \theta = \sqrt{p}$

$$\Rightarrow |\psi\rangle = \cos \theta |B\rangle + \sin \theta |G\rangle$$

$$(U_\psi U_S)^m |\psi\rangle = \cos(2m+1)\theta |B\rangle + \sin(2m+1)\theta |G\rangle$$

U_S acts like reflection across $|B\rangle$

Note: m -times application of $U_\psi U_S$ as in "phase estimation" - blp. But: Cost in general $O(Pm)$

7) $|\psi\rangle$ rotates with rate 2θ in $|B\rangle - |G\rangle$ plane. Use QFT to obtain rate.

Measure last $n+1$ qubits in $\{|G\rangle, |B\rangle\}$ -Basis to obtain

$$\frac{1}{\sqrt{P}} \sum_{m=0}^{P-1} |m\rangle \otimes \left[\cos(2m+1)\theta |B\rangle + \sin(2m+1)\theta |G\rangle \right]$$

$$\frac{1}{C} \sum_{m=0}^{P-1} \underbrace{\sin[(2m+1)\theta]}_{x_m} |m\rangle \quad \left(\begin{array}{c} \text{or} \\ \cos(2m+1)\theta \end{array} \right)$$

where $C = \frac{1}{P} \sum_{m=0}^{P-1} \sin^2[(2m+1)\theta]^2$.

→ Apply QFT to obtain [assume $\theta = \frac{\pi \theta_0}{P}, \theta_0 \in \mathbb{Z}$]

$$\frac{1}{C} \sum_{m=0}^{P-1} x_m |m\rangle, \text{ where } x_m \text{ is given}$$

by IFT of \hat{x}_m , i.e.

$$x_m = \frac{1}{\sqrt{P}} \sum_{k=0}^{P-1} \hat{x}_k e^{\frac{2k\pi i m}{P}}$$

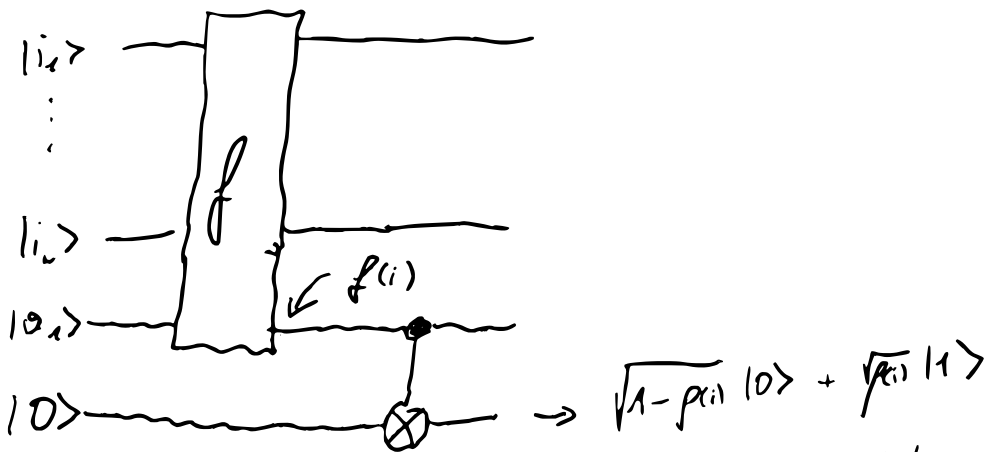
$$= \frac{1}{P} \sum_{k=0}^{P-1} \frac{1}{2i} \left(e^{i(2k+1)\frac{\theta_0}{P}} - e^{-i(2k+1)\frac{\theta_0}{P}} \right) e^{\frac{2k\pi i m}{P}}$$

$$= e^{i\frac{\theta_0}{P}} \frac{1}{2i} \begin{cases} \sqrt{P} & m = \theta_0 \\ 0 & \text{else} \end{cases}$$

8) Final measurement produces $|m\rangle = |\theta_0\rangle$
 and hence $\bar{f} = \sin\left(\frac{\pi\theta_0}{p}\right)^2$

Remark: implementation of Q_f unclear.

if $f: \{0,1\}^n \rightarrow \{0,1\}$, we can use controlled Not-gate to implement



General integrals can always be split into

$$\frac{1}{N} \sum_{i=0}^{N-1} f(i) = \frac{1}{N} \sum_{i=0}^{N-1} \sum_{k=1}^R 2^{-k} f_k(i)$$

i.e., R integration problems for precision 2^{-R} .